

426(4) : Hamilton Jacobi Equation at the Classical Level

In order to apply the HJE to a theory we first consider the generalised Hamiltonian:

$$H = \frac{p_r^2}{2m} + \frac{L^2}{2mr^2} - \frac{n\hbar G}{r} \quad - (1)$$

In the Hamilton Jacobi theory:

$$p_r = \partial S_r / \partial r \quad - (2)$$

and

$$L = \partial S_\phi / \partial \phi \quad - (3)$$

The total action is:

$$S = W(r) + W(\phi) - Et \quad - (4)$$

where E is the total energy and where $W(r)$ and $W(\phi)$ are Hamilton's characteristic functions for the system.

The Hamilton Jacobi equation is therefore: - (5)

$$\frac{1}{2m} \left(\frac{\partial S_r}{\partial r} \right)^2 + \frac{1}{2mr^2} \left(\frac{\partial S_\phi}{\partial \phi} \right)^2 + \frac{\partial S_0}{\partial t} - \frac{n\hbar G}{r} = 0$$

which the actions S_0 , S_r and S_ϕ are separable.

Now write

$$\frac{\partial S_0}{\partial t} = -E \quad - (6)$$

$$\frac{\partial S_\phi}{\partial \phi} = L \quad - (7)$$

and

where E is the total energy and L the angular momentum.

The total energy is the Hamiltonian:

$$H = E \quad (8)$$

and is a constant of motion:

$$\frac{dH}{dt} = 0 \quad (9)$$

Eqn. (9) is the first Ehrenfest equation. The angular momentum (7) is also a constant of motion, so:

$$\frac{dL}{dt} = 0 \quad (10)$$

this is the second Ehrenfest equation. Eq. (5)

therefore becomes:

$$H = \frac{1}{2m} \left(\frac{dS_r}{dr} \right)^2 + U(r) + \frac{L^2}{2mr^2} = E \quad (11)$$

If the interaction potential is:

$$U(r) = -\frac{e^2}{4\pi\epsilon_0 r} \quad (12)$$

Eq. (11) can be transformed into:

$$\frac{d^2 P}{dr^2} + \frac{2m}{\hbar^2} V_{\text{eff}} P = -\frac{2mE}{\hbar^2} P \quad (13)$$

using the quantization conditions:

$$\left(\frac{dS_r}{dr} \right)^2 = -\hbar^2 \frac{d^2 \psi}{dr^2} \quad (14)$$

and

$$L^2 = l(l+1)\hbar^2 \quad (15)$$

3) so the Hamilton Jacobi equation (11) becomes the Schrödinger equation:

$$H\psi = -\frac{\hbar^2}{2m} \nabla^2 \psi - \frac{e^2}{4\pi\epsilon_0 r} \psi = E\psi \quad (16)$$

where

$$\psi(r, \theta, \phi) = R(r) Y(\theta, \phi) \quad (17)$$

and

$$P(r) = rR(r) \quad (18)$$

With these definitions the Hamilton Jacobi equation gives the energy levels of the H atom.

The original Hamilton equations:

$$\dot{p}_r = -\frac{\partial H}{\partial r} \quad (19)$$

and

$$\dot{L} = -\frac{\partial H}{\partial \phi} \quad (20)$$

have been transformed to:

$$\dot{p}_r = -\frac{\partial H}{\partial a_r} = 0 \quad (21)$$

$$\dot{L} = -\frac{\partial H}{\partial \phi} = 0 \quad (22)$$

$$H + \frac{\partial S}{\partial t} = 0 \quad (23)$$

$$p = \frac{\partial S}{\partial q} \quad (24)$$

†) The constants of motion are given by eqs. (6) and (7), which lead to the Evans Eckardt equations:

$$\frac{dH}{dt} = 0 \quad - (25)$$

and

$$\frac{dL}{dt} = 0 \quad - (26)$$

So:

$$\frac{d}{dt} \left(\frac{\partial S_0}{\partial t} \right) = 0 \quad - (27)$$

and

$$\frac{d}{dt} \left(\frac{\partial S_\phi}{\partial \phi} \right) = 0 \quad - (28)$$

but

$$\frac{d}{dt} \left(\frac{\partial S_r}{\partial r} \right) \neq 0 \quad - (29)$$

The action S_r can be found by integrating eq. (11) using a computer, or by recognizing that S_r is related to the associated Legendre functions, the radial wave functions of the H atom. The action S_ϕ is found by integrating:

$$\frac{\partial S_\phi}{\partial \phi} = L = n r^2 \dot{\phi} \quad - (30)$$

a computer. The action S_0 is found by integrating:

$$\frac{\partial S_0}{\partial t} = -E \quad - (31)$$

where for example E are the energy levels of the H atom

5) The quantum of action is h , the reduced Planck constant, so action has the same units as angular momentum. Action is deeply fundamental to all of classical physics, the Hamilton Principle of Least Action being:

$$\int_{t_1}^{t_2} L dt = 0 \quad - (32)$$

where δ is general:

$$\delta S = \int \delta L dt \quad - (33)$$

In order to extend these well known principles to special relativity and a theory, Σ , (1) must be generalized, and the constants of motion and actions defined. Finally, we look for new information on a theory.
