

84(4): Complete Solution of the ECE2 Field

Equations for Planar Orbits

Conditions for derivation of Newtonian orbit:

$$g = -\nabla \Phi + \underline{\omega} \Phi = -\frac{\partial Q}{\partial t} - \omega_0 Q \quad (1)$$

by symmetry. Assume that:

$$\frac{\partial Q}{\partial t} = 0 \quad (2)$$

for a static g . Then:

$$g = -\nabla \Phi + \underline{\omega} \Phi = -\omega_0 Q \quad (3)$$

which has the solution of the previous notes:

$$\omega_0 = -\frac{c}{r} \quad (4)$$

$$Q = -\frac{mb}{c} \frac{r}{r^2} \quad (5)$$

and

$$g = -mb \frac{r}{r^3} \quad (6)$$

Now assume that ω_0 is universal and is the same for all types of planar orbit. Each orbit is therefore characterised by Q , which is different for each orbit.

For a planar orbit:

$$\underline{r} = x \underline{i} + y \underline{j} \quad (7)$$

So :

$$Q_x = -\frac{mG}{c} \frac{x}{x^2+y^2} \quad - (8)$$

$$Q_y = -\frac{mG}{c} \frac{y}{x^2+y^2} \quad - (9)$$

By antisymmetry:

$$\left(\frac{d}{dt} - \omega_y \right) Q_z = - \left(\frac{d}{dt} - \omega_z \right) Q_y \quad - (10)$$

$$\left(\frac{d}{dt} - \omega_z \right) Q_x = - \left(\frac{d}{dt} - \omega_x \right) Q_z \quad - (11)$$

$$\left(\frac{d}{dt} - \omega_x \right) Q_y = - \left(\frac{d}{dt} - \omega_y \right) Q_x \quad - (12)$$

For a planar orbit:

$$\left(\frac{d}{dt} - \omega_x \right) Q_y = - \left(\frac{d}{dt} - \omega_y \right) Q_x \quad - (13)$$

and Eqs. (10) and (11) each reduce to zero on both sides.

The antisymmetry law for planar orbits is

therefore:

$$\boxed{\frac{\partial Q_y}{\partial x} + \frac{\partial Q_x}{\partial y} = \omega_x Q_y + \omega_y Q_x} \quad - (14)$$

In the absence of a gravitomagnetic field:

$$\underline{Q} = \underline{\nabla} \times \underline{Q} - \underline{\omega} \times \underline{Q} = \underline{0} \quad - (15)$$

This means that:

$$\frac{\partial Q_2}{\partial t} - \frac{\partial Q_1}{\partial z} = \omega_1 Q_2 - \omega_2 Q_1 \quad - (16)$$

$$\frac{\partial Q_x}{\partial z} - \frac{\partial Q_z}{\partial x} = \omega_z Q_x - \omega_x Q_z \quad - (17)$$

$$\frac{\partial Q_y}{\partial x} - \frac{\partial Q_x}{\partial y} = \omega_x Q_y - \omega_y Q_x \quad - (18)$$

For a parallel orbit both sides of eqs. (16) and (17) reduce to zero, so in the absence of a gravitational field:

$$\boxed{\frac{\partial Q_y}{\partial x} - \frac{\partial Q_x}{\partial y} = \omega_x Q_y - \omega_y Q_x} \quad - (19)$$

From eqs. (14) and (19):

$$\frac{\partial Q_y}{\partial x} = \omega_x Q_y \quad - (20)$$

$$\frac{\partial Q_x}{\partial y} = \omega_y Q_x \quad - (21)$$

From eqs. (8) and (9):

$$\frac{\partial Q_y}{\partial x} = \frac{mg}{c} \frac{2xy}{(x^2+y^2)^2} \quad - (22)$$

$$\frac{\partial Q_x}{\partial y} = \frac{mg}{c} \frac{2yx}{(x^2+y^2)^2} \quad - (23)$$

so

$$\omega_x = - \frac{2x}{(x^2+y^2)} \quad - (24)$$

$$\omega_y = - \frac{2y}{(x^2+y^2)} \quad - (25)$$

4) and

$$\underline{\omega} = -\frac{2\underline{r}}{r^3} \quad - (26)$$

It is again assumed that this is an intrinsic property of each planet orbit. Therefore the spin connection for vector is

$$\omega^{\mu} = \left(\frac{\omega_0}{c}, \underline{\omega} \right) = - \left(\frac{1}{r}, \frac{2\underline{r}}{r^3} \right) \quad - (27)$$

for Newtonian orbit.

Finally the scalar potential is calculated

$$\begin{aligned} \underline{g} &= -\underline{\nabla} \underline{\Phi} + \underline{\omega} \underline{\Phi} \\ &= -\underline{\nabla} \underline{\Phi} - 2\underline{\Phi} \frac{\underline{r}}{r^3} \\ &= -mG \frac{\underline{r}}{r^3} \end{aligned} \quad - (28)$$

It follows that:

$$\underline{\Phi} = \underline{mG} \quad - (29)$$

This solution obeys the antisymmetry laws (1) and (14).

Note carefully that the sign of $\underline{\Phi}$ is opposite to that of the standard model, i.e.

which has no spin correction for Newtonian gravitation. In Eddington gravitation the spin correction for ^{Newtonian} planar orbits is the same eq. (27).

In this solution:

$$\nabla \times (\omega_0 \underline{a}) = \underline{0} \quad - (30)$$

and

$$\nabla \cdot (\omega_0 \underline{a}) = -4\pi G \rho_m \quad - (31)$$

as in previous notes.

For retrograde precession:

$$\begin{aligned} \underline{g} &= -\frac{MG}{r^3} \frac{\underline{r}}{r^3} = -\omega_0 \underline{a} \quad - (32) \\ &= \frac{c}{r} \underline{a} \end{aligned}$$

so

$$\underline{a} = -\frac{MG}{r^3 c} \frac{\underline{r}}{r^2} \quad - (33)$$

and this obeys eqs. (14) and (19) provided that the spin correction for retrograde precession is defined by eqs. (20) and (21).

So retrograde precession obeys antisymmetry.

For forward precession:

$$\underline{g} = \frac{MG}{r^3} \left(\frac{\underline{r}(\underline{r} \cdot \underline{r})}{c^2} - \underline{r} \right) = -\omega_0 \underline{a}$$

$$= \frac{c}{r} Q \quad - (34)$$

So forward precession is the absence of a gravitomagnetic field obeys antisymmetry provided that the spin connection is defined by eqs. (20) and (21).

So forward precession obeys antisymmetry.

It has been assumed that the scalar spin

connection:

$$\omega_0 = -\frac{c}{r} \quad - (35)$$

is the same for all three types of orbit. Its modulus is a fundamental angular frequency of spacetime:

$$\Omega_0 = 2\pi |\omega_0| = 2\pi \frac{c}{r} \quad - (36)$$

This is the angular frequency of the vacuum particle:

$$E = \hbar \Omega_0 = \gamma mc^2 \quad - (37)$$