

## 238(a): Relativistic Dynamics of a Whirlpool Galaxy

In note 238(8) it was found that:

$$r = \left( 1 - \left( \frac{L_0}{m r_0 c} \right)^2 \left( 1 + \left( \frac{r_0}{r} \right)^2 \right) \right)^{1/2} \frac{L_0}{m r_0} t \quad - (1)$$

This can be solved for  $r^2$  in terms of  $t^2$  as follows. First note that:

$$r^2 = \left( 1 - \left( \frac{L_0}{m r_0 c} \right)^2 \left( \frac{r^2 + r_0^2}{r^2} \right) \right) \left( \frac{L_0}{m r_0} \right)^2 t^2 \quad - (2)$$

$$\text{i.e. } r^4 = \left( r^2 - \left( \frac{L_0}{m r_0 c} \right)^2 (r^2 + r_0^2) \right) \left( \frac{L_0}{m r_0} \right)^2 t^2$$

$$= r^2 \left( 1 - \left( \frac{L_0}{m r_0 c} \right)^2 \right) t^2 \left( \frac{L_0}{m r_0} \right)^2$$

$$- \left( \frac{L_0}{m r_0 c} \right)^2 \left( \frac{L_0}{m r_0} \right)^2 r_0^2 t^2 \quad - (3)$$

$$\text{i.e. } r^4 = A r^2 t^2 - B t^2 \quad - (4)$$

$$\text{where } A = \left( 1 - \left( \frac{L_0}{m r_0 c} \right)^2 \right) \left( \frac{L_0}{m r_0} \right)^2, \quad - (5)$$

$$B = \left( \frac{L_0}{m r_0 c} \right)^2 \left( \frac{L_0}{m r_0} \right)^2 r_0^2 \quad - (6)$$

so for eq. (4):

$$r^2 = \frac{1}{2} \left( At^2 \pm \left( A^2 t^4 - 4Bt^2 \right)^{1/2} \right) \quad - (7)$$

$$\text{i.e. } r^2 = \frac{1}{2} \left( At^2 \pm t \left( A^2 t^2 - 4B \right)^{1/2} \right) \quad - (8)$$

Finally:

$$r = \frac{1}{\sqrt{2}} \left( At^2 \pm t \left( A^2 t^2 - 4B \right)^{1/2} \right)^{1/2}$$

This eq. gives the dependence of  $r$  on  $t$ . In the non-relativistic limit it reduces to:

$$r = \left( \frac{L_0}{m_0} \right) t \quad - (10)$$

if the positive root is used, so:

$$r = \frac{t}{\sqrt{2}} \left( At + \left( A^2 t^2 - 4B \right)^{1/2} \right)^{1/2} \quad - (11)$$

In the non-relativistic limit:

$$A \rightarrow \left( \frac{1}{m_0} \right)^2, \quad B \rightarrow 0. \quad - (12)$$