

Proof of Eq (1b) of Note 231(1) on Commutativity  
of Triple Matrix Products.

To prove

$$(AB)C = A(BC) \quad - (1)$$

for matrices.

Proof

Consider:

$$\begin{aligned} & \left( \begin{bmatrix} A_1 & A_2 \\ A_3 & A_4 \end{bmatrix} \begin{bmatrix} B_1 & B_2 \\ B_3 & B_4 \end{bmatrix} \right) \begin{bmatrix} C_1 & C_2 \\ C_3 & C_4 \end{bmatrix} \\ &= \begin{bmatrix} A_1 B_1 + A_2 B_3 & A_1 B_2 + A_2 B_4 \\ A_3 B_1 + A_4 B_3 & A_3 B_2 + A_4 B_4 \end{bmatrix} \begin{bmatrix} C_1 & C_2 \\ C_3 & C_4 \end{bmatrix} \\ &= \begin{bmatrix} A_1 B_1 C_1 + A_2 B_3 C_1 + A_1 B_2 C_3 + A_2 B_4 C_3 & \dots \\ \dots & \dots \end{bmatrix} \end{aligned} \quad - (2)$$

Similarly consider:

$$\begin{aligned} & \begin{bmatrix} A_1 & A_2 \\ A_3 & A_4 \end{bmatrix} \left( \begin{bmatrix} B_1 & B_2 \\ B_3 & B_4 \end{bmatrix} \begin{bmatrix} C_1 & C_2 \\ C_3 & C_4 \end{bmatrix} \right) \\ &= \begin{bmatrix} A_1 & A_2 \\ A_3 & A_4 \end{bmatrix} \begin{bmatrix} B_1 C_1 + B_2 C_3 & B_1 C_2 + B_2 C_4 \\ B_3 C_1 + B_4 C_3 & B_3 C_2 + B_4 C_4 \end{bmatrix} \quad - (3) \\ &= \begin{bmatrix} A_1 B_1 C_1 + A_1 B_2 C_3 + A_2 B_3 C_1 + A_2 B_4 C_3 & \dots \\ \dots & \dots \end{bmatrix} \end{aligned}$$

It is seen that the products are the same, **QED**  
The same matrix method is used in tensorial

2) algebra. For example:

$$\nabla^a = g_{\mu}^a \nabla^{\mu} \quad - (4)$$

which is two transverse directions is:

$$\begin{bmatrix} \nabla^{(1)} \\ \nabla^{(2)} \end{bmatrix} = \begin{bmatrix} g_{11}^{(1)} & g_{21}^{(1)} \\ g_{11}^{(2)} & g_{21}^{(2)} \end{bmatrix} \begin{bmatrix} \nabla^1 \\ \nabla^2 \end{bmatrix} \quad - (5)$$

$$= \begin{bmatrix} g_{11}^{(1)} \nabla^1 + g_{21}^{(1)} \nabla^2 \\ g_{11}^{(2)} \nabla^1 + g_{21}^{(2)} \nabla^2 \end{bmatrix}$$

It is seen that there is summation over  $\mu$  indices. This is equivalent to matrix algebra. (to repeated)

Next consider:

$$g^{ab} = g_{\mu}^a g^{\mu b} = \begin{bmatrix} g_{11}^{(1)} & g_{21}^{(1)} \\ g_{11}^{(2)} & g_{21}^{(2)} \end{bmatrix} \begin{bmatrix} g^{11} & g^{12} \\ g^{21} & g^{22} \end{bmatrix} \quad - (6)$$

$$= \begin{bmatrix} g_{11}^{(1)} g^{11} + g_{21}^{(1)} g^{21} & g_{11}^{(1)} g^{12} + g_{21}^{(1)} g^{22} \\ g_{11}^{(2)} g^{11} + g_{21}^{(2)} g^{21} & g_{11}^{(2)} g^{12} + g_{21}^{(2)} g^{22} \end{bmatrix}$$

It is seen that there is summation over the repeated  $\mu$  index.

Next consider:

$$g^{ab} g_{\mu\nu} = g_{\mu}^a g^{\mu\nu} g_{\nu b} \quad - (7)$$



3)

To Prove

$$(g_{\mu\nu}^a g^{\mu\nu}) g_{\mu\nu} = g_{\mu\nu}^a (g^{\mu\nu} g_{\mu\nu}) \quad \text{--- (8)}$$

Proof

The left hand side is:

$$\begin{aligned} \text{LHS} &= \begin{bmatrix} g_{11}^{(1)} g'' + g_{21}^{(1)} g^{21} & g_{11}^{(1)} g^{12} + g_{21}^{(1)} g^{22} \\ g_{11}^{(2)} g'' + g_{21}^{(2)} g^{21} & g_{11}^{(2)} g^{12} + g_{21}^{(2)} g^{22} \end{bmatrix} \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix} \\ &= \begin{bmatrix} g_{11}^{(1)} g'' g_{11} + g_{21}^{(1)} g^{21} g_{11} + g_{11}^{(1)} g^{12} g_{21} + g_{21}^{(1)} g^{22} g_{21} & \dots \\ \dots & \dots \end{bmatrix} \quad \text{--- (9)} \end{aligned}$$

The right hand side is:

$$\begin{aligned} \text{RHS} &= \begin{bmatrix} g_{11}^{(1)} & g_{21}^{(1)} \\ g_{11}^{(2)} & g_{21}^{(2)} \end{bmatrix} \left( \begin{bmatrix} g'' & g^{12} \\ g^{21} & g^{22} \end{bmatrix} \begin{bmatrix} g_{11} & g_{12} \\ g_{21} & g_{22} \end{bmatrix} \right) \\ &= \begin{bmatrix} g_{11}^{(1)} & g_{21}^{(1)} \\ g_{11}^{(2)} & g_{21}^{(2)} \end{bmatrix} \begin{bmatrix} g'' g_{11} + g^{12} g_{21} & g'' g_{12} + g^{12} g_{22} \\ g^{21} g_{11} + g^{22} g_{21} & g^{21} g_{12} + g^{22} g_{22} \end{bmatrix} \\ &= \begin{bmatrix} g_{11}^{(1)} g'' g_{11} + g_{11}^{(1)} g^{12} g_{21} + g_{21}^{(1)} g^{21} g_{11} + g_{21}^{(1)} g^{22} g_{21} & \dots \\ \dots & \dots \end{bmatrix} \quad \text{--- (10)} \end{aligned}$$

4) It is seen that the first element of the LHS and RHS products are the same, Q.E.D.

Eq. (15) of note 231(i) is:

$$g^{a\nu} = \nu_{\mu}^a g^{\mu\nu} \quad - (11)$$

indicating that:

$$\boxed{\nu_{\mu}^a = g_{\mu}^a} \quad - (12)$$

This is a fundamental result indicating that the tetrad is a mixed index metric.

In eq. (11)  $g^{\mu\nu}$  is the inverse metric in any mathematical space of any dimension. The metric is  $g_{\mu\nu}$ . So:

$$g^{\mu\nu} g_{\mu\nu} = 1 \quad - (13)$$

Multiply both sides of eq. (11) by  $g_{\mu\nu}$ :

$$g^{a\nu} g_{\mu\nu} = \nu_{\mu}^a g^{\mu\nu} g_{\mu\nu} \quad - (14)$$

$$= \nu_{\mu}^a (g^{\mu\nu} g_{\mu\nu})$$

$$= \nu_{\mu}^a$$

5) So:

$$\boxed{g_{\mu}^{\alpha} = g_{\mu}^{\alpha} = g^{\alpha\nu} g_{\mu\nu}} \quad (15)$$

which is eq. (16) of note 231(1), QED.

The metric  $g^{\alpha\nu}$  is  $g_{\mu}^{\alpha}$  with raised index, and in tensor algebra in any space of any dimension the metric  $g_{\mu\nu}$  or inverse metric  $g^{\mu\nu}$  is always used to lower or raise indices.

A particularly useful example of eq. (15) is when  $\mu$  and  $\nu$  label the Minkowski spacetime. For

$$g_{\mu\nu} = g^{\mu\nu} = \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \quad (16)$$

then:

$$g_{\mu}^{\alpha} = g_{\mu}^{\alpha} = g^{\alpha\nu} \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{bmatrix} \quad (17)$$

Let the space be labeled by  $a = (0), (1), (2), (3)$  — (18)



then:

$$g^{a\nu} = \begin{bmatrix} g^{(0)0} & g^{(0)1} & g^{(0)2} & g^{(0)3} \\ g^{(1)0} & g^{(1)1} & g^{(1)2} & g^{(1)3} \\ g^{(2)0} & g^{(2)1} & g^{(2)2} & g^{(2)3} \\ g^{(3)0} & g^{(3)1} & g^{(3)2} & g^{(3)3} \end{bmatrix} \quad (19)$$

$$\text{and } g_{\mu}^a = g_{\mu}^a = \begin{bmatrix} g^{(0)}_0 & g^{(0)}_1 & g^{(0)}_2 & g^{(0)}_3 \\ g^{(1)}_0 & g^{(1)}_1 & g^{(1)}_2 & g^{(1)}_3 \\ g^{(2)}_0 & g^{(2)}_1 & g^{(2)}_2 & g^{(2)}_3 \\ g^{(3)}_0 & g^{(3)}_1 & g^{(3)}_2 & g^{(3)}_3 \end{bmatrix} \quad (20)$$

$$= \begin{bmatrix} g^{(0)0} & -g^{(0)1} & -g^{(0)2} & -g^{(0)3} \\ g^{(1)0} & -g^{(1)1} & -g^{(1)2} & -g^{(1)3} \\ g^{(2)0} & -g^{(2)1} & -g^{(2)2} & -g^{(2)3} \\ g^{(3)0} & -g^{(3)1} & -g^{(3)2} & -g^{(3)3} \end{bmatrix} \quad (21)$$

In general:

$$g^{a\nu} = h^a h^{\nu} \underline{e}^a \cdot \underline{e}^{\nu} \quad (22)$$

where

$$h^a = |\underline{e}^a| \quad (23)$$

$$h^{\nu} = |\underline{e}^{\nu}| \quad (24)$$

7) So:

$$\eta_{\mu}^{\alpha} = g_{\mu}^{\alpha} = h_{\mu}^{\alpha} \underline{e}^{\alpha} \cdot \underline{e}_{\mu} \quad - (25)$$

where:  $h^{\alpha} = |\underline{e}^{\alpha}|, h_{\mu} = |\underline{e}_{\mu}|. \quad - (26)$

In Minkowski spacetime:

$$\underline{e}^{\alpha} = (\underline{1}, \underline{i}, \underline{j}, \underline{k}) \quad - (27)$$

In  $\mathbb{R}^4$  spacetime labelled by  $(\alpha)$ :

$$\underline{e}^{\alpha} = (\underline{1}, \underline{e}^{(1)}, \underline{e}^{(2)}, \underline{e}^{(3)}) \quad - (28)$$

Therefore  $\underline{e}_{\mu} = (\underline{1}, -\underline{i}, -\underline{j}, -\underline{k}) \quad - (29)$

In  $\mathbb{R}^4$  Minkowski spacetime:

$$|\underline{i}| = |\underline{j}| = |\underline{k}| = 1, \quad - (30)$$

$$\text{so } h^0 = h^1 = h^2 = h^3 = 1 \quad - (31)$$

If the  $\alpha$  basis is chosen with unit scaling factors then:

$$\eta_{\mu}^{\alpha} = g_{\mu}^{\alpha} = \underline{e}^{\alpha} \cdot \underline{e}_{\mu} \quad - (32)$$

This is a fundamental definition that may be used throughout differential geometry.

8) Similarly: 
$$g_{\mu\nu} = \underline{e}^{\mu} \cdot \underline{e}^{\nu} \quad - (33)$$

### The Electromagnetic Potential

This was defined in the early development of ECE theory as:

$$A_{\mu}^a = A_0 g_{\mu}^a \quad - (34)$$

so  $A_{\mu}^a = A_0 g_{\mu}^a = A_0 \underline{e}^a \cdot \underline{e}_{\mu}$

$$A_{\mu}^a = \underline{A}^a \cdot \underline{e}_{\mu} \quad - (35)$$

where  $\underline{A}^a = A_0 \underline{e}^a, \quad - (36)$

i.e.  $A^{(0)} = A_0 \quad - (37)$

$\underline{A}^{(i)} = A_0 \underline{e}^{(i)}, \quad i=1, 2, 3 \quad - (38)$

This definition was introduced originally to deal with circularly polarized radiation, in which

$$\underline{e}^{(1)} = \frac{1}{\sqrt{2}} (\underline{i} - i\underline{j}) e^{i\phi} \quad - (39)$$

$$\underline{e}^{(2)} = \frac{1}{\sqrt{2}} (\underline{i} + i\underline{j}) e^{-i\phi} \quad - (40)$$

$$\phi = \omega t - \kappa z. \quad - (41)$$



9) Here  $\omega$  is angular frequency at instant  $t$  and  $\underline{k}$  is the wave vector at locality  $\underline{r}$  along the  $z$  axis. So the a frame is rotating and translating. w.r.t respect to the  $\mu$  frame.

Equally well, eq. (35) can be written

as:

$$A_{\mu}^a = A_0 g_{\mu}^a = \underline{e}^{(a)} \cdot \underline{A}_{\mu} \quad - (42)$$

This equation makes it clear that the  $e/m$  potential at the ETE level is defined from the potential at the  $u(i)$  level by a dot product with the unit vector  $\underline{e}^{(a)}$ .

For example:

$$A_0^{(0)} = e^{(0)} A_0 = A_0 \quad - (43)$$

$$A_1^{(1)} = \underline{e}^{(1)} \cdot \underline{A}_1 \\ = -\underline{e}^{(1)} \cdot A_0 \underline{i}, \quad - (44)$$

$$A^{(1)1} = A_0 \underline{e}^{(1)} \cdot \underline{i} = A_x^{(1)}$$

$$A^{(1)2} = A_0 \underline{e}^{(1)} \cdot \underline{j} = A_y^{(1)}$$

$$A^{(3)3} = A_0 \underline{e}^{(3)} \cdot \underline{k} = A_z^{(3)}$$

- (45)

10)

In general:

$$\underline{A}^{\alpha\mu} = A \cdot \underline{g}^{\alpha\mu} = \underline{e}^{(\alpha)} \cdot \underline{A}^\mu \quad - (46)$$

where:

$$\underline{e}^{(\alpha)} = (1, \underline{e}^{(1)}, \underline{e}^{(2)}, \underline{e}^{(3)}) \quad - (47)$$

$$\underline{A}^\mu = (A_0, A_0 \underline{i}, A_0 \underline{j}, A_0 \underline{k}) \quad - (48)$$

and by definition:

$$\underline{e}^{(0)} \cdot \underline{A}^0 = A_0 \quad - (49)$$

More generally:

$$\underline{A}^\mu = (A_0, A_x \underline{i}, A_y \underline{j}, A_z \underline{k}) \quad - (50)$$

### Advances over Concepts by Elie Cartan

Cartan introduced the concept of tetrad to define a tangent Minkowski spacetime at point P to a base manifold. His purpose was to define spinors in the general mathematical space. FCF then applied the Cartan method to define the electromagnetic potential in circular polarization. This became a concept of general relativity, i.e. of a propagating frame labelled a spinning and translating with respect to a Minkowski frame  $\mu$ .



11) In this note, eqs. (42) and (46) make it clear that ECE contains more information than  $u(i)$ , it is introduced by forming a metric out of  $e^{(a)}$  and  $A^{\mu}$ . In  $u(i)$  the space is not considered at all, and only one representation used of spacetime. By defining  $A^{\mu}$  as a mixed index metric, the structure equations of Maurer and Cartan can be used to define torsion and curvature. On the  $u(i)$  level these concepts do not exist in electrodynamics. The tetrad postulate can be used to find the spin connection. The first structure equation gives the relation between field and potential, and the Cartan identity between torsion and curvature gives the field equations of electrodynamics. The connection becomes an intrinsic part of electrodynamics and a generally covariant unified field theory emerges.

Eqs. (42) and (46) show that all the dynamic information is contained in the moving frame with unit vector  $e^{(a)}$ .

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