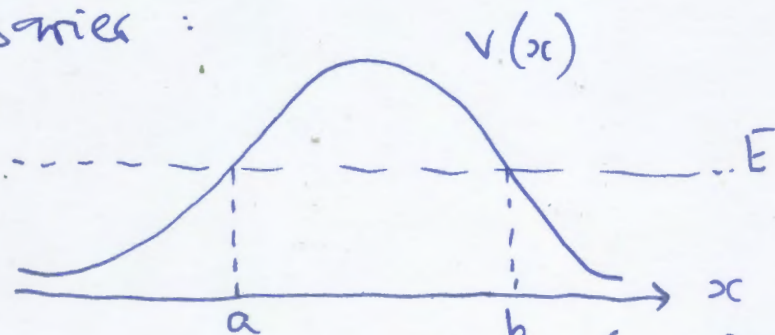


229(1) : Tunneling Through an Arbitrary Barrier.

Consider the series :

Fig(1)



Using the Wentzel, Kramers, Brillouin (WKB) approx.,
 then as in MerzSader pp. 126 ff:

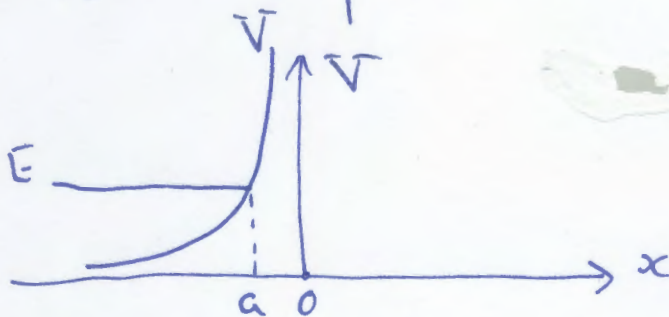
$$T = 4 / \left(2\theta + \frac{1}{2\theta} \right)^2 \quad - (1)$$

where:

$$\theta = \exp \left(\int_a^b \kappa(x) dx \right) \quad - (2)$$

Consider a nucleus of charge $Z_1 e$ fusing
 with a nucleus of charge $Z_2 e$. In order to do this
 the Coulomb barrier has to be quantum tunneled:

Fig(2)



Then:

$$V = \frac{Z_1 Z_2 e^2}{4\pi \epsilon_0 x} \quad - (3)$$

for

$$x < 0. \quad - (4)$$

2) The part a is defined by:

$$E = \frac{Z_1 Z_2 e^2}{4\pi\epsilon_0 a} \quad - (5)$$

and it is assumed that

$$b = 0 \quad - (6)$$

Then:

$$K = \left(\frac{2m}{\hbar}\right)^{1/2} (V - E)^{1/2} \quad - (7)$$

$$\int_a^0 K dx = \frac{(2mE)^{1/2}}{\hbar} \int_a^0 \left(\frac{a}{x} - 1\right)^{1/2} dx \quad - (8)$$

for $x < a$ - (9)

Therefore

$$\theta = \exp\left(\frac{(2mE)^{1/2}}{\hbar} \int_a^0 \left(\frac{a}{x} - 1\right)^{1/2} dx\right) \quad - (10)$$

and

$$T = \frac{4}{\left(\frac{\partial\theta}{\partial a} + \frac{1}{\partial\theta}\right)^2} \quad - (11)$$

$$E = Z_1 Z_2 e^2 / a \quad - (12)$$

Therefore T depends on the mass m of the incoming α particle, on Z_1 , Z_2 and a as follows:

$$A = \exp \left(\frac{e}{f} \left(\frac{2m z_1 z_2}{4\pi \epsilon_0 a} \right)^{1/2} \int_a^{\infty} \left(\frac{a}{x} - 1 \right)^{1/2} dx \right) \quad - (13)$$

Here:

$$e = 1.60219 \times 10^{-19} \text{ C}$$

$$f = 6.62618 \times 10^{-34} \text{ Js}$$

$$4\pi \epsilon_0 = 1.11265 \times 10^{-10} \text{ J}^{-1} \text{ C}^2 \text{ m}^{-1}$$

Therefore:

$$A = \exp \left(\left(\frac{e^2 m z_1 z_2}{2\pi f^2 \epsilon_0 a} \right)^{1/2} \int_a^{\infty} \left(\frac{a}{x} - 1 \right)^{1/2} dx \right) \quad - (14)$$

Units Check

$$\left(\frac{e^2 m z_1 z_2}{2\pi f^2 \epsilon_0 a} \right)^{1/2} = \left(\frac{\text{C}^2 \text{ kg m}}{\text{J}^2 \text{ s}^2 \text{ J}^{-1} \text{ C}^2 \text{ m}^{-1} \text{ m}} \right)^{1/2} = \left(\frac{\text{kg m}}{\text{J s}^2} \right)^{1/2}$$

$$\text{J} = \text{kg m}^2 \text{ s}^{-2}$$

so the exponent is dimensionless as required
 because the units of the integral are m. ✓✓

Now define: $A = e^{x_1} \quad - (15)$

then $T = \frac{4}{\left(2e^{x_1} + \frac{1}{2} e^{-x_1} \right)^2} \quad - (16)$

4) where:

$$x_1 = \left(\frac{e^2 m Z_1 Z_2}{2\pi \hbar^2 \epsilon_0 a} \right)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx \quad - (17)$$

$$x_1 = \left(\frac{e^2}{2\pi \hbar^2 \epsilon_0} \right)^{1/2} \left(\frac{m Z_1 Z_2}{a} \right)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx \quad - (18)$$

The fine structure constant is:

$$\alpha = \frac{e^2}{4\pi \hbar c \epsilon_0} \quad - (19)$$

$$\text{So } x_1 = \left(\frac{2\alpha c}{\hbar} \right)^{1/2} \left(\frac{m Z_1 Z_2}{a} \right)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx$$

-(20)

So:

where m is the mass of the proton.

$$x_1 = \left(\frac{2\alpha c m}{\hbar} \right)^{1/2} \left(\frac{Z_1 Z_2}{a} \right)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx \quad - (21)$$

where

$$\alpha = 0.007297351 \quad - (22)$$

$$c = 2.997925 \times 10^8 \text{ ms}^{-1}$$

$$m = 1.67265 \times 10^{-27} \text{ kg}$$

$$\hbar = 1.05459 \times 10^{-34} \text{ Js}$$

Therefore:

$$x_1 = 8.33 \times 10^6 \left(\frac{z_1 z_2}{a} \right)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx$$

—(23)

Here

$$E = \frac{z_1 z_2}{4\pi \epsilon_0 a^2}, \quad \text{—(24)}$$

So

$$\frac{z_1 z_2}{a} = 4\pi \epsilon_0 a E, \quad \text{—(25)}$$

and

$$x_1 = 8.33 \times 10^6 (4\pi \epsilon_0 a E)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx$$

$$x_1 = 87.9 (aE)^{1/2} \int_a^0 \left(\frac{a}{x} - 1 \right)^{1/2} dx$$

—(26)

and

$$T = \frac{4}{\left(2e^{x_1} + \frac{1}{2} e^{-x_1} \right)^2} \quad \text{—(27)}$$

For small a and E , T becomes large, and fusion is possible.