

181(2) : Hamiltonian Jacobi Equation for Fermi Equation.

Define the ECE wave equation as:

$$(\square + R_0 + R_1) \phi = 0 \quad - (1)$$

$$R = R_0 + R_1, \quad - (2)$$

$$R_0 = \left( \frac{m_0 c}{\hbar} \right)^2, \quad - (3)$$

$$R_1 = \hbar^\mu \hbar_\mu = \frac{\omega^2}{c^2} - \kappa^2 \quad - (4)$$

The advantage of this format is that it is based on a constant Hamiltonian:

$$H = \frac{1}{2} m_0 c^2 \quad - (5)$$

The interaction of a field/particle with particle of observed mass  $m_0$  is described by  $R_1$ .

In the absence of interaction:

$$\left( \square + \left( \frac{m_0 c}{\hbar} \right)^2 \right) \phi = 0 \quad - (6)$$

whose classical description is:

$$p^\mu p_\mu = m_0^2 c^2 \quad - (7)$$

Eq. (7) can be factorized into the free fermion equations:

$$(E - c \underline{\sigma} \cdot \underline{p}) \phi^R = m_0 c^2 \phi^L \quad - (8)$$

$$(E + c \underline{\sigma} \cdot \underline{p}) \phi^L = m_0 c^2 \phi^R \quad - (9)$$

2) In the position representation  $\underline{p}$  is the vector valued relativistic momentum:

$$\underline{p} = \gamma m \underline{v} \quad - (10)$$

and  $E$  is the total relativistic energy:

$$E = \gamma m_0 c^2 \quad - (11)$$

Multiply (8) by (9) to obtain:

$$E^2 - c^2 \underline{p}^2 = m_0^2 c^4 \quad - (12)$$

which is eq. (7).

The interaction is described by:

$$p^\mu \rightarrow p^\mu - \not{R} \kappa^\mu \quad - (13)$$

where

$$R_1 = \kappa^\mu \kappa_\mu = \frac{\omega^2}{c^2} - \kappa^2 \quad - (14)$$

Here:

$$p^\mu = \left( \frac{E}{c}, \underline{p} \right), \quad \kappa^\mu = \left( \frac{\omega}{c}, \underline{\kappa} \right) \quad - (15)$$

so

$$E \rightarrow E - \not{R} \omega, \quad \underline{p} \rightarrow \underline{p} - \not{R} \underline{\kappa} \quad - (16)$$

Therefore eqs. (8) and (9) become:

$$\left( E - \not{R} \omega - c \underline{\sigma} \cdot \left( \underline{p} - \not{R} \underline{\kappa} \right) \right) \phi^R = m_0 c^2 \phi^L \quad - (17)$$

$$\left( E - \not{R} \omega + c \underline{\sigma} \cdot \left( \underline{p} - \not{R} \underline{\kappa} \right) \right) \phi^L = m_0 c^2 \phi^R \quad - (18)$$

Multiply eqs. (17) and (18) to obtain:

3) From eq. (18):

$$\phi^R = \frac{1}{m_0 c^2} (E - \hbar \omega + c \underline{\sigma} \cdot (\underline{p} - \hbar \underline{k})) \phi^L \quad (19)$$

Therefore eq. (17) becomes:

$$(E - \hbar \omega - c \underline{\sigma} \cdot (\underline{p} - \hbar \underline{k})) ((E - \hbar \omega + c \underline{\sigma} \cdot (\underline{p} - \hbar \underline{k})) \phi^L) = m_0^2 c^4 \phi^L \quad (20)$$

If eqs. (17) and (18) are multiplied together, the result is

$$(E - \hbar \omega)^2 - c^2 (\underline{p} - \hbar \underline{k}) \cdot (\underline{p} - \hbar \underline{k}) = m_0^2 c^4 \quad (21)$$

which is the relativistic Hamilton-Jacobi equation:

$$\boxed{(\underline{p}^\mu - \hbar \underline{k}^\mu)(\underline{p}_\mu - \hbar \underline{k}_\mu) = m_0^2 c^2} \quad (22)$$

We will incorporate  $\hbar \underline{k}$  into eq. (22).

This is done by expanding the left hand side of the equation:

$$\underline{p}^\mu \underline{p}_\mu + \hbar^2 \underline{k}^\mu \underline{k}_\mu - \hbar (\underline{k}^\mu \underline{p}_\mu + \underline{p}^\mu \underline{k}_\mu) = m_0^2 c^2 \quad (23)$$

Write the left hand side as:

$$\underline{p}^\mu \underline{p}_\mu - \hbar^2 R_1 = m_0^2 c^2 \quad (24)$$

where  $R_1$  is defined by eq. (1). Eq. (24) is the classical equivalent of eq. (1) with:

4) 
$$p^\mu = i\hbar \partial^\mu \quad - (25)$$

so 
$$p^\mu p_\mu = -\hbar^2 \partial^\mu \partial_\mu = -\hbar^2 \square \quad - (26)$$

Therefore: 
$$\hbar^2 R_1 = -\hbar^2 \kappa^\mu \kappa_\mu + \hbar (\kappa^\mu p_\mu + p^\mu \kappa_\mu) \quad - (27)$$

Finally use: 
$$p^\mu \rightarrow p^\mu + \hbar \kappa^\mu \quad - (28)$$

which is the minimal prescription for describing the interaction of the free particle with another field/particle. Therefore in eq. (27)

$$p_\mu = \hbar \kappa_\mu \quad - (29)$$

and 
$$R_1 = \kappa^\mu \kappa_\mu \quad - (30)$$

which is eq. (4), QED.

It has been shown that eq. (1) is the quantized version of the Hamilton-Jacobi equation (22), given the minimal prescription (28). Particle collisions can be described by:

$$R = R_0 + R_1 \quad - (31)$$

$$= \left(\frac{m_0 c}{\hbar}\right)^2 + \kappa^\mu \kappa_\mu$$

5) i.e.

$$R = \left( \frac{m_0 c}{\hbar} \right)^2 + \frac{\omega^2}{c^2} - k^2 \quad - (32)$$

Here  $\omega$  and  $k$  are the angular frequency and wavenumbers associated with the matter wave interacting with the free particle of rest mass  $m_0$ .

The wave equation for the interaction is:

$$\left( \square + \frac{\omega^2}{c^2} - k^2 + \left( \frac{m_0 c}{\hbar} \right)^2 \right) \psi = 0 \quad - (33)$$

The correct description of particle interaction at the classical level is therefore:

$$P_{\mu}^{\mu} = \hbar^2 k_{\mu} k_{\mu} + m_0^2 c^2 \quad - (34)$$

i.e.

$$\frac{P_{\mu}^{\mu}}{\hbar^2} = R_1 + R_0 \quad - (35)$$

or

$$\boxed{P_{\mu}^{\mu} = \hbar^2 (R_1 + R_0)} \quad - (36)$$

Eq. (36) may now be applied to particle collisions.